

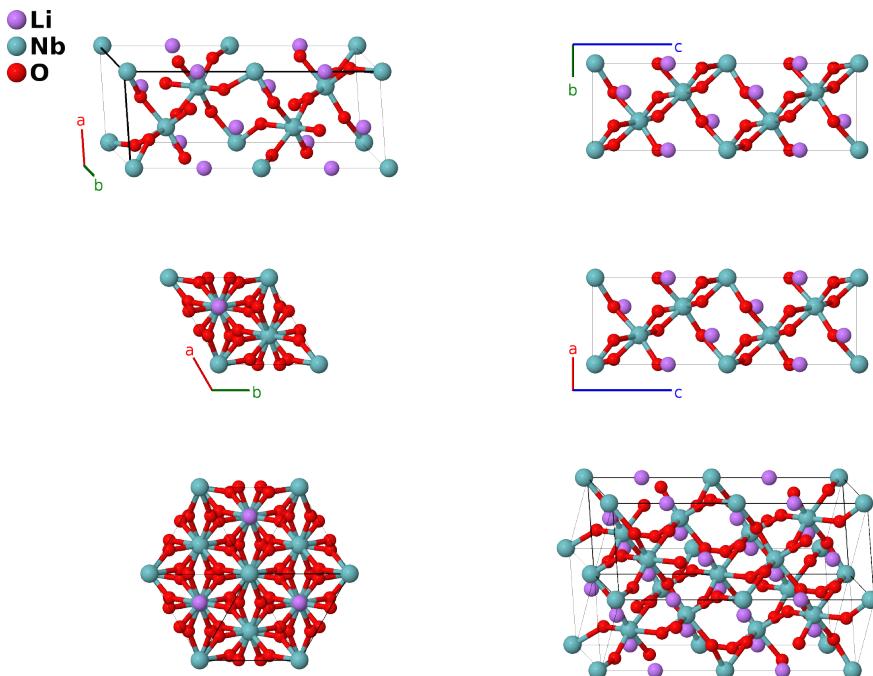
# Ferroelectric LiNbO<sub>3</sub> Structure: ABC3\_hR10\_161\_a\_a\_b-001

This structure originally had the label ABC3\_hR10\_161\_a\_a\_b. Calls to that address will be redirected here.

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<https://aflow.org/p/FZF0>

[https://aflow.org/p/ABC3\\_hR10\\_161\\_a\\_a\\_b-001](https://aflow.org/p/ABC3_hR10_161_a_a_b-001)



**Prototype** LiNbO<sub>3</sub>

**AFLOW prototype label** ABC3\_hR10\_161\_a\_a\_b-001

**ICSD** 81243

**Pearson symbol** hR10

**Space group number** 161

**Space group symbol** *R*3c

**AFLOW prototype command**

```
aflow --proto=ABC3_hR10_161_a_a_b-001  
--params=a,c/a,x1,x2,x3,y3,z3
```

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## Other compounds with this structure

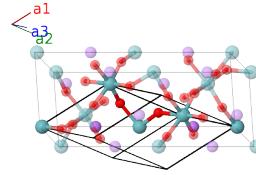
BiFeO<sub>3</sub>, CsPbF<sub>3</sub>

- 
- This is the ferroelectric phase of LiNbO<sub>3</sub>, which is stable below 1430K. There is also a high-temperature paraelectric phase.

- In the special case  $c/a = \sqrt{6}$ ,  $z_1 = 1/4$ ,  $z_2 = 0$ ,  $x_3 = 1/2$ ,  $y_3 = 0$ ,  $z_3 = 0$  this reduces to a double unit cell version of the cubic perovskite ( $E2_1$ ) structure. This sets the angle between the rhombohedral primitive vectors to  $60^\circ$ . Experimentally the value is about  $56^\circ$ .

## Rhombohedral primitive vectors

$$\begin{aligned}\mathbf{a}_1 &= \frac{1}{2}a\hat{\mathbf{x}} - \frac{\sqrt{3}}{6}a\hat{\mathbf{y}} + \frac{1}{3}c\hat{\mathbf{z}} \\ \mathbf{a}_2 &= \frac{1}{\sqrt{3}}a\hat{\mathbf{y}} + \frac{1}{3}c\hat{\mathbf{z}} \\ \mathbf{a}_3 &= -\frac{1}{2}a\hat{\mathbf{x}} - \frac{\sqrt{3}}{6}a\hat{\mathbf{y}} + \frac{1}{3}c\hat{\mathbf{z}}\end{aligned}$$



## Basis vectors

	Lattice coordinates	Cartesian coordinates	Wyckoff position	Atom type
$\mathbf{B}_1$	$x_1 \mathbf{a}_1 + x_1 \mathbf{a}_2 + x_1 \mathbf{a}_3$	$cx_1 \hat{\mathbf{z}}$	(2a)	Li I
$\mathbf{B}_2$	$(x_1 + \frac{1}{2}) \mathbf{a}_1 + (x_1 + \frac{1}{2}) \mathbf{a}_2 + (x_1 + \frac{1}{2}) \mathbf{a}_3$	$c(x_1 + \frac{1}{2}) \hat{\mathbf{z}}$	(2a)	Li I
$\mathbf{B}_3$	$x_2 \mathbf{a}_1 + x_2 \mathbf{a}_2 + x_2 \mathbf{a}_3$	$cx_2 \hat{\mathbf{z}}$	(2a)	Nb I
$\mathbf{B}_4$	$(x_2 + \frac{1}{2}) \mathbf{a}_1 + (x_2 + \frac{1}{2}) \mathbf{a}_2 + (x_2 + \frac{1}{2}) \mathbf{a}_3$	$c(x_2 + \frac{1}{2}) \hat{\mathbf{z}}$	(2a)	Nb I
$\mathbf{B}_5$	$x_3 \mathbf{a}_1 + y_3 \mathbf{a}_2 + z_3 \mathbf{a}_3$	$\frac{1}{2}a(x_3 - z_3) \hat{\mathbf{x}} - \frac{\sqrt{3}}{6}a(x_3 - 2y_3 + z_3) \hat{\mathbf{y}} + \frac{1}{3}c(x_3 + y_3 + z_3) \hat{\mathbf{z}}$	(6b)	O I
$\mathbf{B}_6$	$z_3 \mathbf{a}_1 + x_3 \mathbf{a}_2 + y_3 \mathbf{a}_3$	$-\frac{1}{2}a(y_3 - z_3) \hat{\mathbf{x}} + \frac{\sqrt{3}}{6}a(2x_3 - y_3 - z_3) \hat{\mathbf{y}} + \frac{1}{3}c(x_3 + y_3 + z_3) \hat{\mathbf{z}}$	(6b)	O I
$\mathbf{B}_7$	$y_3 \mathbf{a}_1 + z_3 \mathbf{a}_2 + x_3 \mathbf{a}_3$	$-\frac{1}{2}a(x_3 - y_3) \hat{\mathbf{x}} - \frac{\sqrt{3}}{6}a(x_3 + y_3 - 2z_3) \hat{\mathbf{y}} + \frac{1}{3}c(x_3 + y_3 + z_3) \hat{\mathbf{z}}$	(6b)	O I
$\mathbf{B}_8$	$(z_3 + \frac{1}{2}) \mathbf{a}_1 + (y_3 + \frac{1}{2}) \mathbf{a}_2 + (x_3 + \frac{1}{2}) \mathbf{a}_3$	$-\frac{1}{2}a(x_3 - z_3) \hat{\mathbf{x}} - \frac{\sqrt{3}}{6}a(x_3 - 2y_3 + z_3) \hat{\mathbf{y}} + \frac{1}{6}c(2x_3 + 2y_3 + 2z_3 + 3) \hat{\mathbf{z}}$	(6b)	O I
$\mathbf{B}_9$	$(y_3 + \frac{1}{2}) \mathbf{a}_1 + (x_3 + \frac{1}{2}) \mathbf{a}_2 + (z_3 + \frac{1}{2}) \mathbf{a}_3$	$\frac{1}{2}a(y_3 - z_3) \hat{\mathbf{x}} + \frac{\sqrt{3}}{6}a(2x_3 - y_3 - z_3) \hat{\mathbf{y}} + \frac{1}{6}c(2x_3 + 2y_3 + 2z_3 + 3) \hat{\mathbf{z}}$	(6b)	O I
$\mathbf{B}_{10}$	$(x_3 + \frac{1}{2}) \mathbf{a}_1 + (z_3 + \frac{1}{2}) \mathbf{a}_2 + (y_3 + \frac{1}{2}) \mathbf{a}_3$	$\frac{1}{2}a(x_3 - y_3) \hat{\mathbf{x}} - \frac{\sqrt{3}}{6}a(x_3 + y_3 - 2z_3) \hat{\mathbf{y}} + \frac{1}{6}c(2x_3 + 2y_3 + 2z_3 + 3) \hat{\mathbf{z}}$	(6b)	O I

## References

- [1] H. Boysen and F. Altorfer, *A neutron powder investigation of the high-temperature structure and phase transition in LiNbO<sub>3</sub>*, Acta Crystallogr. Sect. B **50**, 405–414 (1994), doi:10.1107/S0108768193012820.