Fe₈N $(D2_g)$ Structure: A8B_tI18_139_deh_a-001

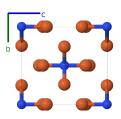
This structure originally had the label ASB_tI18_139_deh_a. Calls to that address will be redirected here.

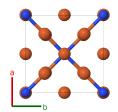
Cite this page as: D. Hicks, M. J. Mehl, M. Esters, C. Oses, O. Levy, G. L. W. Hart, C. Toher, and S. Curtarolo, *The AFLOW Library of Crystallographic Prototypes: Part 3*, Comput. Mater. Sci. **199**, 110450 (2021), doi: 10.1016/j.commatsci.2021.110450.

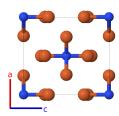
https://aflow.org/p/6VR8

https://aflow.org/p/A8B_tI18_139_deh_a-001









Prototype Fe_8N

AFLOW prototype label A8B_tI18_139_deh_a-001

Strukturbericht designation $D2_q$

ICSD 654562

Pearson symbol tI18

Space group number 139

Space group symbol I4/mmm

AFLOW prototype command aflow --proto=A8B_tI18_139_deh_a-001

--params= $a, c/a, z_3, x_4$

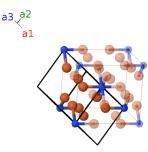
- (Jack, 1951) and others refer to this structure as α'' -Fe₁₆N₂.
- (Yamashita, 2012) determined the structure by examining crystals containing various amounts of α'' -Fe₁₆N₂ and α -Fe (bcc iron). We use their data from the sample that was 90% α'' -Fe₁₆N₂. Their measurements were confirmed by first principles calculations and are in agreement with the first principles calculations of (Sims, 2012).
- The ICSD entry is from (Jack, 1951).

Body-centered Tetragonal primitive vectors

$${f a_1} = -rac{1}{2}a\,{f \hat x} + rac{1}{2}a\,{f \hat y} + rac{1}{2}c\,{f \hat z}$$

$$\mathbf{a_2} \quad = \quad \frac{1}{2}a\,\mathbf{\hat{x}} - \frac{1}{2}a\,\mathbf{\hat{y}} + \frac{1}{2}c\,\mathbf{\hat{z}}$$

$$\mathbf{a_3} = \frac{1}{2}a\,\mathbf{\hat{x}} + \frac{1}{2}a\,\mathbf{\hat{y}} - \frac{1}{2}c\,\mathbf{\hat{z}}$$



Basis vectors

		Lattice coordinates		Cartesian coordinates	Wyckoff position	Atom type
$\mathbf{B_1}$	=	0	=	0	(2a)	NΙ
$\mathbf{B_2}$	=	$rac{3}{4}{f a}_1 + rac{1}{4}{f a}_2 + rac{1}{2}{f a}_3$	=	$rac{1}{2}a\mathbf{\hat{y}}+rac{1}{4}c\mathbf{\hat{z}}$	(4d)	Fe I
$\mathbf{B_3}$	=	$\frac{1}{4}\mathbf{a}_1 + \frac{3}{4}\mathbf{a}_2 + \frac{1}{2}\mathbf{a}_3$	=	$\frac{1}{2}a\hat{\mathbf{x}} + \frac{1}{4}c\hat{\mathbf{z}}$	(4d)	Fe I
$\mathbf{B_4}$	=	$z_3\mathbf{a}_1+z_3\mathbf{a}_2$	=	$cz_3\mathbf{\hat{z}}$	(4e)	Fe II
${f B_5}$	=	$-z_3\mathbf{a}_1-z_3\mathbf{a}_2$	=	$-cz_3\mathbf{\hat{z}}$	(4e)	Fe II
${f B_6}$	=	$x_4 \mathbf{a}_1 + x_4 \mathbf{a}_2 + 2x_4 \mathbf{a}_3$	=	$ax_4\mathbf{\hat{x}} + ax_4\mathbf{\hat{y}}$	(8h)	Fe III
$\mathbf{B_7}$	=	$-x_4\mathbf{a}_1 - x_4\mathbf{a}_2 - 2x_4\mathbf{a}_3$	=	$-ax_4\mathbf{\hat{x}}-ax_4\mathbf{\hat{y}}$	(8h)	Fe III
${f B_8}$	=	$x_4 \mathbf{a}_1 - x_4 \mathbf{a}_2$	=	$-ax_4\mathbf{\hat{x}} + ax_4\mathbf{\hat{y}}$	(8h)	Fe III
$\mathbf{B_9}$	=	$-x_4\mathbf{a}_1+x_4\mathbf{a}_2$	=	$ax_4\mathbf{\hat{x}}-ax_4\mathbf{\hat{y}}$	(8h)	Fe III

References

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- [3] H. Sims, W. H. Butler, M. Richter, K. Koepernik, E. Şaşioğlu, C. Friedrich, and S. Blügel, *Theoretical investigation into the possibility of very large moments in Fe*₁₆N₂, Phys. Rev. B **86**, 174422 (2012), doi:10.1103/PhysRevB.86.174422.