

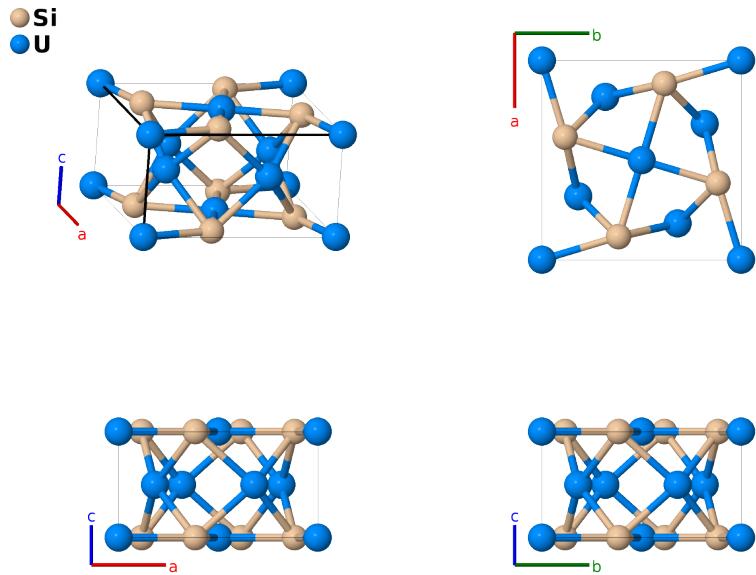
Si_2U_3 ($D5_a$) Structure: A2B3_tP10_127_g_ah-001

This structure originally had the label A2B3_tP10_127_g_ah. Calls to that address will be redirected here.

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<https://aflow.org/p/6QGS>

https://aflow.org/p/A2B3_tP10_127_g_ah-001



| | |
|------------------------------------|---|
| Prototype | Si_2U_3 |
| AFLOW prototype label | A2B3_tP10_127_g_ah-001 |
| Strukturbericht designation | $D5_a$ |
| ICSD | 73695 |
| Pearson symbol | tP10 |
| Space group number | 127 |
| Space group symbol | $P4/mbm$ |
| AFLOW prototype command | <code>aflow --proto=A2B3_tP10_127_g_ah-001 --params=a, c/a, x₂, x₃</code> |

Other compounds with this structure

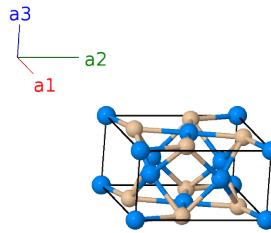
Al₂Th₃, B₂Mo₃, B₂Nb₃, B₂Ta₃, B₂V₃, Be₂Nb₃, Be₂Ta₃, Ga₂Nb₃, Ga₂Ta₃, Ga₂Zr₃, Ge₂Th₃, Pd₂Dy₃, Pd₂Ho₃, Si₂Hf₃, Si₂Th₃, Si₂Zr₃

- If we consider the Si₂ dimers as a pseudo-atom then this is a tetragonal distortion of the Cu₃Au ($L1_2$) structure.
- Mo₂FeB₂ is the ternary form of this structure.

- We use the lattice constants taken by (Remschnig, 1992) at 1200°C, while the ICSD entry uses what are apparently the room temperature values.

Simple Tetragonal primitive vectors

$$\begin{aligned}\mathbf{a}_1 &= a \hat{\mathbf{x}} \\ \mathbf{a}_2 &= a \hat{\mathbf{y}} \\ \mathbf{a}_3 &= c \hat{\mathbf{z}}\end{aligned}$$



Basis vectors

| | Lattice coordinates | = | Cartesian coordinates | Wyckoff position | Atom type |
|-------------------|---|---|---|------------------|-----------|
| \mathbf{B}_1 | 0 | = | 0 | (2a) | U I |
| \mathbf{B}_2 | $\frac{1}{2} \mathbf{a}_1 + \frac{1}{2} \mathbf{a}_2$ | = | $\frac{1}{2} a \hat{\mathbf{x}} + \frac{1}{2} a \hat{\mathbf{y}}$ | (2a) | U I |
| \mathbf{B}_3 | $x_2 \mathbf{a}_1 + (x_2 + \frac{1}{2}) \mathbf{a}_2$ | = | $a x_2 \hat{\mathbf{x}} + a (x_2 + \frac{1}{2}) \hat{\mathbf{y}}$ | (4g) | Si I |
| \mathbf{B}_4 | $-x_2 \mathbf{a}_1 - (x_2 - \frac{1}{2}) \mathbf{a}_2$ | = | $-a x_2 \hat{\mathbf{x}} - a (x_2 - \frac{1}{2}) \hat{\mathbf{y}}$ | (4g) | Si I |
| \mathbf{B}_5 | $-(x_2 - \frac{1}{2}) \mathbf{a}_1 + x_2 \mathbf{a}_2$ | = | $-a (x_2 - \frac{1}{2}) \hat{\mathbf{x}} + a x_2 \hat{\mathbf{y}}$ | (4g) | Si I |
| \mathbf{B}_6 | $(x_2 + \frac{1}{2}) \mathbf{a}_1 - x_2 \mathbf{a}_2$ | = | $a (x_2 + \frac{1}{2}) \hat{\mathbf{x}} - a x_2 \hat{\mathbf{y}}$ | (4g) | Si I |
| \mathbf{B}_7 | $x_3 \mathbf{a}_1 + (x_3 + \frac{1}{2}) \mathbf{a}_2 + \frac{1}{2} \mathbf{a}_3$ | = | $a x_3 \hat{\mathbf{x}} + a (x_3 + \frac{1}{2}) \hat{\mathbf{y}} + \frac{1}{2} c \hat{\mathbf{z}}$ | (4h) | U II |
| \mathbf{B}_8 | $-x_3 \mathbf{a}_1 - (x_3 - \frac{1}{2}) \mathbf{a}_2 + \frac{1}{2} \mathbf{a}_3$ | = | $-a x_3 \hat{\mathbf{x}} - a (x_3 - \frac{1}{2}) \hat{\mathbf{y}} + \frac{1}{2} c \hat{\mathbf{z}}$ | (4h) | U II |
| \mathbf{B}_9 | $-(x_3 - \frac{1}{2}) \mathbf{a}_1 + x_3 \mathbf{a}_2 + \frac{1}{2} \mathbf{a}_3$ | = | $-a (x_3 - \frac{1}{2}) \hat{\mathbf{x}} + a x_3 \hat{\mathbf{y}} + \frac{1}{2} c \hat{\mathbf{z}}$ | (4h) | U II |
| \mathbf{B}_{10} | $(x_3 + \frac{1}{2}) \mathbf{a}_1 - x_3 \mathbf{a}_2 + \frac{1}{2} \mathbf{a}_3$ | = | $a (x_3 + \frac{1}{2}) \hat{\mathbf{x}} - a x_3 \hat{\mathbf{y}} + \frac{1}{2} c \hat{\mathbf{z}}$ | (4h) | U II |

References

[1] K. Remschnig, T. L. Bihan, H. Noël, and P. Rogl, *Structural chemistry and magnetic behavior of binary uranium silicides*, J. Solid State Chem. **97**, 391–399 (1992), doi:10.1016/0022-4596(92)90048-Z.

Found in

[1] P. Villars, K. Cenzual, R. Gladyshevskii, O. Shcherban, V. Dubenskyy, V. Kuprysyuk, I. Savesyuk, and R. Zaremba, *Landolt-Börnstein - Group III Condensed Matter* (2012). URL http://materials.springer.com/lb/docs/sm_lbs_978-3-642-22847-6_359.